

Transition to oscillatory behavior and breakup of shock profiles in a model of relativistic fluid dynamics

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Dissipative relativistic fluid dynamics plays an important role in high energy physics and modern cosmology. Recently new descriptions for it have been proposed. A particularly promising approach for modelling pure radiation fluids has been suggested by Bemfica, Disconzi and Noronha in [1]. This model is given by the four field PDE formulation

$$(1) \quad \partial_\beta(T^{\alpha\beta} + \Delta T^{\alpha\beta}) = 0$$

with ideal part $T^{\alpha\beta}$ of the energy momentum tensor and dissipative part

$$\Delta T^{\alpha\beta} = -\eta B_0^{\alpha\beta} + \mu B_1^{\alpha\beta} + \nu B_2^{\alpha\beta},$$

where $\eta > 0$ quantifies the viscosity and $\mu, \nu > 0$ are “causalizing” parameters. This ansatz augments the Eckart description [6] Sec. 2.11, which corresponds to the choice $\mu = \nu = 0$. As it was shown in [1],[3] there is a range $\mathcal{C} \subset \mathbb{R}_+^3$ of parameters (η, μ, ν) that make this model causal.

For an ideal plane shock wave, i.e. a discontinuous solution to the Euler equations $T^{\alpha\beta} = 0$ of the form

$$\psi(x) = \begin{cases} \psi_-, & x^\beta \xi_\beta < 0 \\ \psi_+, & x^\beta \xi_\beta > 0, \end{cases}$$

a traveling wave solution $\psi_B(x^\beta \xi_\beta)$ to (1) that has the same end states, $\psi_B(\pm\infty) = \psi_\pm$, is called its dissipative profile and is regarded as the dissipative counterpart of the ideal shock wave.

These profiles thus are heteroclinic orbits of a planar dynamical system of the form

$$(2) \quad B(\psi_B, \eta, \mu, \nu) \psi'_B = F(\psi_B, q)$$

where $q \in \mathbb{R}^2$ is the parameter that identifies the shock wave. In contrast to the situation in [2], the fact that $B(\cdot, \cdot, \cdot, \cdot)$ is in general not definite and not even always regular leads to a rich variety of phenomena as the parameters vary.

While as shown in [3] the mere existence of shock profiles always fails for some shocks for parameter values (η, μ, ν) in the interior of \mathcal{C} , the present talk focuses on the situation with (η, μ, ν) on the boundary of \mathcal{C} , a case that is referred to as sharply causal since then a characteristic speed of the dissipation tensor equals the speed of light.

We show that in the sharply causal case there are always Lax shocks for which the system (2) has a spiraling focus at ψ_+ , so that a corresponding traveling wave oscillates at that end. This oscillatory behavior could well indicate a lack of dynamical stability, as it was observed in other applications [4],[5]. It is associated with a non standard form of saddle-node bifurcation that we discuss in some detail. For still increased amplitude, profiles cease to exist, as the matrix B is singular on a submanifold of the state space.

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